PFe₃(CO)₁₁ (average Fe-C(bridging), 2.04 and 1.87 Å).⁷ As Dahm and Jacobson pointed out,⁷ this asymmetry need not be the result of crystal packing forces but could be inherent in the bonding of the parent $Fe_3(CO)_{12}$, in which there is some evidence for unsymmetrical bridging carbonyl groups.⁵ The C–O bond lengths average 1.17 Å, and the distances in the bridging carbonyl groups are not significantly longer than the average.

All the other bond lengths and valency angles (Table IV) are quite similar to those in related molecules. The angles in the $Fe_3(CO)_{10}$ moiety are similar to those in $(C_6H_5)_3PFe_3(CO)_{11}$,⁷ and the dimensions of the diarsine ligand are quite close to those in $LFe_2(CO)_{6,8}$ apart from the differences caused by the nonplanarity of the ligand in the latter compound. The valency angles at arsenic show deviations from the exact tetrahedral value, the Fe-As-C angles (113-121°) being larger than the C-As-C angles (99-104°).

The magnitudes of the principal axes of the thermal vibration ellipsoids of the iron and arsenic atoms are given in Table V. The largest vibrations are approximately perpendicular to the plane of the iron triangle.

All the intermolecular distances correspond to nor-

TABLE V Magnitudes (Å, $\sigma = 0.005-0.008$ Å) of the Principal Axes OF THE THERMAL VIBRATION ELLIPSOIDS OF THE IRON AND ARSENIC ATOMS

AND INSERTCHIONS						
				Molecule 2		
	Axis 1	Axis 2	Axis 3	Axis 1	Axis 2	Axis 3
Fe (1)	0.186	0.233	0.235	0.179	0.249	0.292
Fe(2)	0.165	0.175	0.277	0.195	0.223	0.298
Fe(3)	0.190	0.192	0.261	0.190	0.222	0.273
As(4)	0.177	0.179	0.249	0.194	0.212	0.243
As(5)	0.166	0.218	0.260	0.193	0.232	0.240

mal van der Waals interactions, the closest approaches being about 3.1 Å. The molecules are arranged (Figure 2) so that an oxygen atom of a bridging carbonyl group of each molecule in the asymmetric unit is approximately equidistant from three of the four terminal carbonyl groups of the unique iron atom of the other molecule. Another, simplified view of the molecular packing is shown in Figure 3.

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Some Reactions of the Octahalodirhenate(III) Ions. VIII. Definitive Structural Characterization of the Octabromodirhenate(III) Ion¹

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The compound Cs₂Re₂Br₈ has been prepared and its structure investigated by single-crystal X-ray diffractometry. The existence of the $\text{Re}_2\text{Br}_8^{2-}$ ion with virtual D_{4h} symmetry is confirmed. The Re-Re bond has a length of 2.228 (4) Å, very similar to that found in related compounds. The Re-Br bonds have an average length of 2.478 ± 0.006 Å. The set of eight bromine atoms defines a right, square parallelepiped with mean basal edge 3.39 ± 0.02 Å and mean vertical edge 3.48 ± 0.04 Å. Each cesium ion is surrounded by 11 bromine atoms (Cs \cdots Br distances in the range 3.63–4.17 Å), while each Re₂Br₈²⁻⁷ ion has a cesium ion in approximately symmetrical contact with each of its four vertical faces. The space group is Pbca with $a = 12.625 \pm 0.004$ Å, $b = 12.953 \pm 0.004$ Å, and $c = 10.141 \pm 0.003$ Å. For Z = 4, d(calcd) = 5.11 g cm⁻³ ($d(\text{measd}) = 12.625 \pm 0.004$ Å, d(calcd) = 5.11 g cm⁻³ ($d(\text{measd}) = 12.625 \pm 0.004$ Å, d(calcd) = 5.11 g cm⁻³ ($d(\text{measd}) = 12.625 \pm 0.004$ Å, d(calcd) = 5.11 g cm⁻³ ($d(\text{measd}) = 12.625 \pm 0.004$ Å, $d(\text{calcd}) = 12.625 \pm 0.004$ 5.10 g cm⁻³) and the Re₂Br₈²⁻ units are required to have $\overline{1}$ symmetry. The structure was solved using 770 reflections within the sphere bounded by $2\theta = 40^{\circ}$ (Mo K α) collected with a counter diffractometer and refined by full-matrix least-squares methods with anisotropic temperature parameters to a conventional R factor of 7.1% and a weighted R of 5.4%.

Introduction

A number of molecules and ions containing a quadruple Re-Re bond are known² and several have been structurally characterized with high accuracy by X-ray crystallography. Among these are the $\text{Re}_2\text{Cl}_3^{2-}$ ion³⁻⁵ in which the Re–Re distance is 2.241 (7) Å, the Re₂Cl₆- $[P(C_2H_5)_3]_2$ molecule⁶ in which Re-Re = 2.222 (3) Å, and the $(C_3H_7CO_2)_4Re_2(ReO_4)_2$ and $(C_6H_5CO_2)_4Re_2Cl_2$

molecules⁷ in which the Re–Re distances are 2.251 (2) and 2.235 (2) Å, respectively. In $Re_2Cl_4(O_2CCH_3)_2$ - $(H_2O)_2$, where the four oxygen atoms of the bridging acetato groups lie on one vertical face of the Re₂X₈ parallelepiped and the four Cl atoms define the opposite, parallel vertical face, with approximately the same nonbonded distances as in the $Re_2Cl_8^{2-}$ ion itself, the Re-Re distance is reported⁸ to be 2.224 (5) Å. In these species, then, the bond length is relatively in-(7) (a) C. Calvo, N. C. Jayadevan, C. J. L. Lock, and R. Restivo, Can. J. Chem., 48, 219 (1970); (b) M. J. Bennett, W. K. Bratton, F. A. Cotton, and W. R. Robinson, Inorg. Chem., 7, 1570 (1968).

⁽¹⁾ Supported in part by the U. S. Atomic Energy Commission.

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⁽⁸⁾ N. A. Porai-Koshits, presenting the paper of V. G. Kuznetsov, P. A. Koz'min, and M. D. Suraxskaya, Eighth Congress of the International Union of Crystallography, Stony Brook, N. Y., Aug 1969.

variant, as might be expected for a very strong bond under roughly similar sets of steric constraints. In the case of the $(C_6H_5CO_2)_2Re_2I_4$ molecule,⁵ where the coordination about the rhenium atoms differs significantly from that in the other three molecules the Re-Re distance is appreciably shorter, viz., 2.198 (1) Å.

In connection with studies of the vibrational spectra and other properties of species containing quadruply bonded rhenium atoms the structure of the Re₂Br₈²⁻⁻ ion took on special interest. While it was expected to be isostructural with the $Re_2Cl_8^{2-}$ ion, the exact dimensions, especially the length of the Re-Re bond, were required. Two previously reported structures9 only served to increase the uncertainty concerning the exact value. The two substances studied were crystallographically different forms of (C5H6N)2Re2Br8 (incorrectly formulated as $(C_5H_6N)_2H_2Re_2Br_8)$. For one, where the structure was said to have been refined by least squares (to R = 14.2%), the distance was reported as 2.207 ± 0.003 Å while for the other form, which was not refined, the distance was reported as 2.27 Å with no stated uncertainty interval. The first of these distances seemed unreasonably short; the second seemed unreasonably long though of such dubious accuracy that it could not in any case be taken seriously. A new, accurate investigation was considered necessary.

We report here the preparation of $Cs_2Re_2Br_8$, which is a new compound and the first one, to our knowledge, in which an $\operatorname{Re}_2 X_8^{2-}$ ion occurs in a crystal containing only a monatomic counterion and no additional molecules such as H₂O. The virtue of such a substance is that it makes possible the computation of a lattice energy which in turn may be introduced into a thermodynamic cycle involving the energy of the metal-tometal bond. The results of an accurate single crystal study of the structure of this substance are reported here.

Experimental Section

Preparation.—Ten grams of KReO₄, 10 ml of 50% aqueous $\rm H_3PO_2,~and~20~ml$ of $48\%~\rm HBr$ were refluxed for 2 hr under prepurified nitrogen. Then 8.0 g of CsBr was added; the mixture was stirred for a few minutes and allowed to cool. The resulting brown precipitate was filtered off and washed with ethanol. This crude product was stirred with 6 M HBr at $\sim 80^{\circ}$; the brownish green solution was filtered from the much less soluble yellow Cs2ReBr6 and slowly cooled to produce tabular black crystals of Cs₂Re₂Br₈. This recrystallization was repeated to give samples for elemental analysis and the crystallographic examination. Anal. Calcd for Cs2Re2Br8: Re, 29.15; Br, 50.09. Found: Re, 29.2; Br, 49.9.

Collection of X-Ray Data.---A single crystal was examined by Weissenberg and precession photography. It was found to belong to the orthorhombic system with the following systematic absences: $0kl, k = 2n + 1; \ h0l, l = 2n + 1; \ hk0, h = 2n + 1.$ The space group uniquely indicated is Pbca $(D_{2h}^{15}, no. 61)$.¹⁰ The unit cell dimensions were derived from values of 2θ for the h00, 0k0, and 00l sets of reflections which were carefully measured on a manually operated General Electric XRD-5 quarter-circle diffractometer using Mo K α radiation $\lambda(K\alpha_1)$ 0.70926; $\lambda(K\alpha_2)$ 0.71354 Å). The measurements extended to at least the 14th

order on all three axes, and the α doublet was resolved at the higher angles. Measurements on two crystals were in good agreement giving $a = 12.625 \pm 0.004$ Å, $b = 12.953 \pm 0.004$ Å, and $c = 10.141 \pm 0.003$ Å, where the standard deviations are the authors' estimates and exceed the average deviations of the individual measurements. The density calculated with these dimensions and assuming Z = 4 with a formula weight of 1277.48 is 5.11 g cm⁻³. The density measured pycnometrically was 5.10 g cm⁻³. Since the general position in the space group Pbca has a multiplicity of 8, it is required that Re₂Br₈²⁻ units lie on centers of symmetry.

The crystal, glued to the end of a glass fiber, was so aligned that the b axis was coincident with the ϕ axis of the goniometer. This crystal was a rhombic tablet 0.018 mm thick (between principal (100) and $(\overline{100})$ faces) whose periphery was bounded by approximately equally developed {111} faces and smaller {010} faces. The two greatest thicknesses of the crystal, parallel to the b and c axes and to the diagonals of the rhombus, were both 0.06 mm. This morphology is typical of the crystals obtained.

The intensities of 770 reflections within one octant of the sphere bounded by $2\theta = 40^{\circ}$ were measured by a coupled $\theta - 2\theta$ scan technique. The diffractometer was operated with a zirconium foil filter between the crystal and scintillation detector, a pulse height discriminator set to accept $\sim 95\%$ of the Mo K α peak, and a takeoff angle of $\sim 2^{\circ}$ to the molybdenum X-ray tube anode. Each reflection was measured by performing a 1.33 $^{\circ}$ (2θ) scan, centered on the peak, at a rate of $2^{\circ}/\text{min}$, accumulating P counts in 40 sec. Twenty-second background counts, B_1 and B_2 , were taken at each end of the scan range. Four strong reflections which were remeasured at regular intervals showed only random fluctuations in the range $\pm 2\%$. All of the systematic absences in the region of the data set were also measured and found to confirm the extinction rules previously observed on films.

Values for the magnitudes of the observed structure factors, $|F_{o}|$, and their standard deviations, $\sigma(|F_{o}|)$, were calculated from the experimental data as $|F_o| = (I/Lp)^{1/2}$ and $\sigma(|F_o|) = (P + B_1 + B_2 + (0.04I)^2)^{1/2}/2Lp|F_o|$, where $I = P - B_1 - B_2$ (if P $-B_1 - B_2 \le 0$, $I = 0.5(P + B_1 + B_2)^{1/2}$ and Lp is the Lorentzpolarization factor.

Solution and Refinement.¹¹-A three-dimensional Patterson function was calculated. From it and the assumption of 4/mmm $(D_{4\mathrm{h}})$ symmetry in the $\mathrm{Re}_{2}\mathrm{Br}_{8}{}^{2-}$ ion, the positions of all atoms were found. After four cycles of least-squares refinement¹² with all atoms assigned isotropic temperature factors, the residuals $R_1 = \Sigma ||F_0| - |F_0||/\Sigma |F_0|$ and $R_2 = [\Sigma w (|F_0| - |F_0|)^2 / \Sigma w |F_0|^2]^{1/2}$ were 0.14 and 0.12, respectively. Scattering factors employed in these and all subsequent cycles were those of Cromer and Waber¹³ for neutral Cs, Re, and Br atoms and included anomalous dispersion corrections as given by Cromer.14

Before proceeding with the refinement, an absorption correction was applied to the data using the measured dimensions of the crystal and the absorption coefficient of 402 cm^{-1} . Calculated transmission factors ranged from 0.271 to 0.542. The ratio of $|F_{o}|_{\max}^{2}$ to $|F_{o}|_{\min}^{2}$ for the 020 and 040 reflections over $0^{\circ} \leq \phi \leq$ 360° at $\chi = 90^{\circ}$ was 2.6; the absorption correction reduced this to 1.3. No attempt was made to modify the calculation in order to reduce the ratio further. Four more cycles of least-squares refinement on the corrected data in which all atoms were allowed anisotropic temperature factors converged to $R_1 = 0.071$ and R_2 = 0.054. In the last cycle, no parameter shifted by more than 3% of its estimated standard deviation. The standard deviation of an observation of unit weight was 1.13, average $w(F_{\circ}$ -

⁽⁹⁾ P. A. Koz'min, V. G. Kuznetzov, and Z. V. Popova, Zh. Strukt. Khim.,

<sup>6, 651 (1965).
(10) &</sup>quot;International Tables for X-Ray Crystallography," Vol. I, Kynoch Press, Birmingham, England, 1962.

⁽¹¹⁾ Local, modified versions of the following programs were used: (1) MIXG2, D. P. Shoemaker, diffractometer settings; (2) FORDAP, A. Zalkin, Fourier syntheses; (3) GONO9, W. C. Hamilton, absorption corrections; (4) SFLS5, C. T. Prewitt, least-squares refinements; (5) MGEOM, J. S. Wood, and DISTAN, D. P. Shoemaker, interatomic distances and angles; (6) ORTEP, C. Johnson, thermal ellipsoid drawings.

⁽¹²⁾ The function minimized was $\Sigma w \langle |F_0| - |F_c| \rangle^2$, where the weights, w. were the reciprocals of $\sigma^2(|F_0|)$. (13) D. T. Cromer and J. T. Waber, Acta Crystallogr., **18**, 104 (1965).

⁽¹⁴⁾ D. T. Cromer, ibid., 18, 17 (1965).

 TABLE I

 Observed and Calculated Structure Factors for Cs2Re2Bir (in Units of 0.1 Electron)

 $F_{\rm o}$)² values were essentially invariant with respect to $|F_{\rm o}|$ and $\lambda^{-1} \sin \theta$, and the largest peak on a final difference Fourier synthesis was 1.85 e⁻Å⁻³. A list of the observed and final calculated structure factors was examined for evidence of secondary extinction effects (*i.e.*, $(F_{\rm o}/F_{\rm c})$ vs. I). None was found. The list is presented as Table I.

Results

The fractional atomic coordinates and isotropic temperature parameters are listed in Table II. Anisotropic temperature parameters are given in Table III. The thermal vibration ellipsoids are displayed in Figure 1. The dimensions of the $\text{Re}_2\text{Br}_8^{2-}$ ion are given in Table IV. The shorter Cs to Br and intermolecular Br to Br distances are given in Table V, and the crystal packing is shown in Figure 2.

		TABLE II		
Final	POSITIONAL AND	ISOTROPIC	THERMAL I	PARAMETERS ^a
Atom	x	У	z	$B,^b$ Å ²
Re	0.0154(1)	0.0704(1)	0.060	0(1) 2.38
Cs	-0.3072(2)	0.1359(2)	-0.160	8(2) 4.91
Br(1)	0.0640(3)	0.0029(3)	0.280	1(3) 3.84
Br(2)	-0.1585(3)	0.1146(3)	0.155	8(4) 3.64
Br(3)	-0.0155(3)	0.2189(3)	-0.089	9(3) 3.99
Br(4)	0.2068(3)	0.1052(3)	0.024	7(4) 3.52

^a Estimated standard deviations occurring in the least significant figure are given in parentheses. ^b B's are isotropic thermal parameters equivalent to the anisotropic tensors of Table III.

TANER III

			TABLE	111		
Α	NISOTRO	PIC TEM	IPERATUR	e Parami	ETERS ^{a,b} (>	< 10 ⁴)
Atom	β_{11}	B 22	\$ 83	β_{12}	β18	\$ 28
Re	40(1)	29(1)	63(2)	-2(1)	-2(1)	-3(1)
Cs	88(2)	53(2)	135(3)	21(2)	-34(3)	-26(2)
Br(1)	66(3)	63(3)	74(1)	2(3)	-14(3)	5(3)
Br(2)	59(3)	48(3)	96(4)	6(3)	14(3)	8(3)
Br(3)	76(4)	36(3)	115(5)	8(3)	2(3)	18(3)
Br(4)	43(3)	47(3)	112(4)	-5(2)	-3(3)	0(3)

^{*a*} Estimated standard deviations occurring in the least significant figure appear in parentheses. ^{*b*} Anisotropic temperature factors were defined as $\exp[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23})]$.

 $\begin{array}{c} Table \ IV\\ Interatomic \ Distances \ (\AA) \ and \ Angles \ (deg)\\ In \ the \ Re_2 Br_8^{2-} \ Ion^a \end{array}$

Bond Dista	ances	Nonbond Dis	stances
Re–Re	2.228(4)	Br(1)- $Br(2)$	3.402(5)
Re-Br(1)	2.474(4)	Br(2)-Br(3)	3.361(5)
Re-Br(2)	2.469(4)	Br(3)- $Br(4)$	3.377(5)
Re-Br(3)	2.483(4)	Br(4)- $Br(1)$	3.423(5)
Re-Br(4)	2.484(4)	Br(1)– $Br(3)'$	3.51(1)
		Br(2)– $Br(4)'$	3.44(1)
	Bond	Angles	
Re'-Re-Br(1)	104.3(1)	Br(2)-Re- $Br(3)$	85.5(2)
Re'-Re-Br(2)	104.4(1)	Br(3)-Re- $Br(4)$	85.7(2)
Re'Re-Br(3)	105.8(1)	Br(4)-Re- $Br(1)$	87.3(2)
Re'-Re-Br(4)	103.8(1)	Br(1)-Re- $Br(3)$	149.9(2)
Br(1)-Re- $Br(2)$	87.0(2)	Br(2)-Re- $Br(4)$	151.7(2)

^a Estimated standard deviations occurring in the least significant figure appear in parentheses.

	Тав	le V		
J	INTERIONIC DI	STANCES ^a (Å)		
Anion at (0,	0, 0)	Anion at $(0, 0.5, -0.5)$		
Cs-Br(4)'	3.64	Cs-Br(2)	4.17	
Cs-Br(2)	3.73 2.76	Between ani	ons	
Cs-Br(3)	3.90	Br(2)-Br(4) Br(2)-Br(3)	3.66 3.81	
Anion at $(-0.5, 0, -0.5)$		Br(1)- $Br(2)$	3.85	
Cs-Br(4)	3.71	Br(1)-Br(3)	3.97	
Cs-Br(1)'	3.76	Br(1)- $Br(4)$	4.06	
Cs-Br(2)	3.77	Br(3)-Br(4)	4.23	
Cs-Br(3)'	3.80	$B_1(2)$ - $Br(4)$	4.41	
Anion at $(-0,$	5, 0.5, 0)			
Cs-Br(4)	3.63			
Cs-Br(3)	4.11			

 a The estimated standard deviations of these distances are all 0.005 Å or less. All other Cs–Br and Br–Br distances are greater than 4.5 Å.

Discussion

The geometric configuration of the $\text{Re}_2\text{Br}_8^{2-}$ ion is, as expected, that of a nearly cubic, right square parallelepiped. Figure 1 shows the structure with mean di-



Figure 1.—The Re₂Br₈²⁻ ion showing the atom numbering scheme and mean dimensions. Intervals indicated with \pm are root-mean-square deviations from the mean and exceed the standard deviations of individual values used to compute the mean. Primed atoms are related by the center of symmetry to unprimed atoms with the same number.

mensions and the atom numbering scheme. The individual dimensions are given in Table IV. The Re_2 - Br_8^{2-} ion lies at a crystallographic center of symmetry. Its dimensions are such, however, that it is more chemically meaningful to think of it as having virtual D_{4h} symmetry. The eclipsed configuration and short Re–Re distance are diagnostic² of the quadruple Re–Re bond.

The Re-Re bond length, 2.228 (4) Å, is near the middle of the range, roughly 2.22–2.25 Å, covered by other Re-Re quadruple bonds. This is significant since it shows that the quadruple bond distance is rela-



Figure 2.—The Cs₂Re₂Br₈ structure projected down a.

tively constant, 2.235 ± 0.015 Å, provided general stereochemistry remains essentially the same. The increase to 2.29 Å in the case of Re₂Cl₅(CH₃SCH₂CH₂-SCH₃)₂ becomes meaningful and supports the view that the latter compound contains only a triple bond.¹⁵

The crystal packing, shown in Figure 2, is not particularly remarkable. Each $\text{Re}_2\text{Br}_3^{2-}$ ion is surrounded by four Cs⁺ ions, each in approximately symmetrical contact with one of the vertical faces of the right square parallelepiped which it forms. Each Cs⁺ ion is surrounded by 11 bromine atoms, four from each of two anions, two from a basal edge of another, and one corner of a fourth. This array of bromine atoms does not approximate to any particular symmetric polyhedron. Eight of the Cs-Br distances are in the range 3.63-3.80Å, while the other three are appreciably longer, *viz.*, 3.90, 4.11, and 4.17 Å.

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Synthesis of Sodium Cyanotrihydroborate and Sodium Isocyanotrihydroborate

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A convenient synthetic procedure for sodium cyanotrihydroborate (NaBH₃CN) is reported. Sodium isocyanotrihydroborate has also been isolated in mixtures with the normal salt and characterized. The isomerization of NaBH₃NC to Na-BH₃CN has been investigated and is found to be H⁺ and CN⁻ catalyzed. Boron nmr data are reported for BH₃CN⁻, BH₃NC⁻, and BH₃CNBH₃⁻. The infrared spectra of NaBH₃NC, NaBD₃NC, and NaBH₃CNBH₃ have been recorded and the fundamental vibrations assigned and compared with those reported earlier for NaBH₃CN and NaBD₃CN.

Introduction

Wittig and Raff reported the first synthesis of a cyanotrihydroborate in 1951.² They prepared the lithium salt by allowing lithium borohydride to react

(1) (a) PRF Graduate Fellowship 1969-1970; abstracted in part from the Ph.D. thesis of J. R. Berschied, Kansas State University, 1970. (b) To whom inquiries should be addressed. (c) Kansas State University. (d) Ventron Corp.

(2) G. Wittig and P. Raff, Z. Naturforsch. B, 6, 225 (1951).

with excess HCN in diethyl ether at 100° under pressure. These authors noted the unusual hydrolytic stability of the product toward acids. Thereafter, only one brief publication appeared on the reducing power of LiBH₃CN, and that report indicates limited applicability.³

(3) G. Drefahl and E. Keil, J. Prakt. Chem., 6, 80 (1958).

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